

Random repeated quantum interactions and random invariant states

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Repeated quantum interactions and quantum channels

Invariant states and the asymptotic induced ensemble

Random environment states and i.i.d. interaction unitaries

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- ▶ Schrödinger equation

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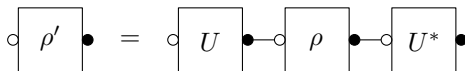
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- ▶ Graphical form



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- ▶ The system $\rho \in \mathcal{M}_d^{1,+}(\mathbb{C})$ is coupled to a (possibly unknown/inaccessible) environment described by a state $\beta \in \mathcal{M}_{d'}^{1,+}(\mathbb{C})$ and undergoes a “closed” unitary dynamics described by a matrix $U \in \mathcal{U}(dd')$

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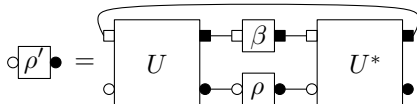
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Remark

The transposition map

$$\begin{aligned} T : \mathcal{M}_d(\mathbb{C}) &\rightarrow \mathcal{M}_d(\mathbb{C}) \\ X &\mapsto X^T \end{aligned}$$

is not completely positive.

Quantum channels

A linear map $\Phi : \mathcal{M}_d(\mathbb{C}) \rightarrow \mathcal{M}_d(\mathbb{C})$ is a quantum channel iff one of the following holds:

1. **Stinespring dilation** There exists a finite dimensional Hilbert space $\mathbb{C}^{d'}$, a density matrix $\beta \in \mathcal{M}_{d'}^{1,+}(\mathbb{C})$ and an unitary operation $U \in \mathcal{U}(dd')$ such that

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2. **Kraus decomposition** There exists an integer k and matrices $L_1, \dots, L_k \in \mathcal{M}_d(\mathbb{C})$ such that

$$\Phi(X) = \sum_{i=1}^k L_i X L_i^*, \quad \forall X \in \mathcal{M}_d(\mathbb{C})$$

and

$$\sum_{i=1}^k L_i^* L_i = I_d.$$

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It can be shown that the dimension of the ancilla space in the Stinespring dilation theorem can be chosen $d' = d^2$ and β can be chosen to be a rank one projector. A similar result holds for the number of Kraus operators: one can always find a decomposition with $k = d^2$ operators.

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The *Choi rank* of a quantum channel Φ is the least positive integer k such that Φ admits a Kraus decomposition with k operators L_i .

Two examples

- ▶ For $U \in \mathcal{U}(d)$, define the **unitary conjugation channel**

$$\Phi_U(X) = UXU^*.$$

One can check that the spectrum of Φ_U is

$$\text{spec}(\Phi_U) = \{\lambda_1 \bar{\lambda}_2 \mid \lambda_1, \lambda_2 \in \text{spec}(U)\}.$$

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- ▶ The **depolarizing channel** $\Phi_{\text{dep}} : \mathcal{M}_d(\mathbb{C}) \rightarrow \mathcal{M}_d(\mathbb{C})$ is given by

$$\Phi_{\text{dep}}(X) = \text{Tr}(X) \frac{I}{d}.$$

It has eigenvalues 1 (with multiplicity 1) and 0 (with multiplicity $d^2 - 1$).

Spectral properties of channels

Proposition

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1. Φ has at least one invariant element which is a density matrix;
2. Φ has trace operator norm 1;
3. Φ has spectral radius 1;
4. Φ satisfies the Schwarz inequality

$$\forall X \in \mathcal{M}_d(\mathbb{C}), \quad \Phi(X)^* \Phi(X) \leq \|\Phi(I)\| \Phi(X^* X).$$

Asymptotic states for a class of channels

Let \mathcal{C} be the set of all quantum channels that have 1 as a simple eigenvalue and all other eigenvalues are contained in the **open** unit disc.

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Proposition

Consider a quantum channel $\Phi \in \mathcal{C}$. Then, for all density matrices $\rho_0 \in \mathcal{M}_d^{1,+}(\mathbb{C})$,

$$\lim_{n \rightarrow \infty} \Phi^n(\rho_0) = \rho_\infty,$$

where ρ_∞ is the unique invariant state of Φ .

A model of random quantum channels

Fix two integers $d, d' \geq 2$ and a density matrix $\beta \in \mathcal{M}_{d'}^{1,+}(\mathbb{C})$. To an unitary matrix $U \in \mathcal{U}(dd')$, associate the channel

$$\Phi^{U,\beta}(X) = \text{Tr}_{d'} [U(X \otimes \beta)U^*].$$

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Choosing U random from the **Haar distribution** $\mathfrak{h}_{dd'}$ on the unitary group, we obtain a quantum channel-valued random variable (β is fixed)

$$\begin{aligned} \mathcal{U}(dd') &\rightarrow \mathcal{L}(\mathcal{M}_d(\mathbb{C})) \\ U &\mapsto \Phi^{U,\beta}. \end{aligned}$$

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Question: What are the properties of a *generic* quantum channel ?

Almost all quantum channels are in \mathcal{C}

Theorem

Let β be a fixed density matrix of size d' . If U is a random unitary matrix distributed along the Haar invariant probability $\mathfrak{h}_{dd'}$ on $\mathcal{U}(dd')$, then $\Phi^{U,\beta} \in \mathcal{C}$ almost surely.

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Corollary

For almost all unitary matrices $U \in \mathcal{U}(dd')$, the channel $\Phi^{U,\beta}$ has an unique invariant state ρ_∞ and for all density matrices ρ_0 ,

$$\lim_{n \rightarrow \infty} \left(\Phi^{U,\beta} \right)^n (\rho_0) = \rho_\infty.$$

Strictly positive and irreducible channels

Definition

A positive map $\Phi : \mathcal{M}_d(\mathbb{C}) \rightarrow \mathcal{M}_d(\mathbb{C})$ is called

- ▶ **strictly positive** (or positivity improving) if $\Phi(X) > 0$ for all $X \geq 0$;

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Proposition

A positive linear map $\Phi : \mathcal{M}_d(\mathbb{C}) \rightarrow \mathcal{M}_d(\mathbb{C})$ is irreducible if and only if the map $(I + \Phi)^{d-1}$ is strictly positive.

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Theorem

If Ψ is a unital, irreducible map on $\mathcal{M}_d(\mathbb{C})$ which satisfies the Schwarz inequality (eg. the dual of an irreducible quantum channel Φ), then the set of peripheral (i.e. modulus one) eigenvalues is a (possibly trivial) subgroup of the unit circle \mathbb{T} . Moreover, every peripheral eigenvalue is simple and the corresponding eigenspaces are spanned by unitary elements of $\mathcal{M}_d(\mathbb{C})$.

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Corollary

The peripheral eigenvalues of an irreducible quantum channel are simple and contained in the finite set

$$\{\xi \in \mathbb{T} \mid \exists 1 \leq n \leq d^2 \text{ s.t. } \xi^n = 1\}.$$

Necessary and sufficient conditions for irreducibility

We denote by $\text{Lat}(T)$ the lattice of invariant subspaces of an operator $T \in \mathcal{M}_d(\mathbb{C})$.

Proposition (Farenick)

Consider a completely positive map $\Phi : \mathcal{M}_d(\mathbb{C}) \rightarrow \mathcal{M}_d(\mathbb{C})$ defined by

$$\Phi(X) = \sum_{i=1}^k L_i X L_i^*,$$

with $L_i \in \mathcal{M}_d(\mathbb{C})$, $i = 1, \dots, k$. Then Φ is irreducible if and only if $\bigcap_{i=1}^k \text{Lat}(L_i)$ is trivial.

Necessary and sufficient conditions for irreducibility

Proposition (the Shemesh criterion)

Two matrices $A, B \in \mathcal{M}_d(\mathbb{C})$ have a common eigenvector if and only if

$$\bigcap_{i,j=1}^{d-1} \ker [A^i, B^j] \neq \{0\}.$$

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More generally, if A and B have a common invariant subspace of dimension k (for $1 \leq k \leq d-1$), then their k -th wedge powers have a common eigenvector, and hence (we put $n = \binom{d}{k}$)

$$\bigcap_{i,j=1}^{n-1} \ker [(A^{\wedge k})^i, (B^{\wedge k})^j] \neq \{0\}.$$

Almost all quantum channels are irreducible

- ▶ Write the matrix U defining a quantum channel Φ as a $d' \times d'$ matrix of blocks in $\mathcal{M}_d(\mathbb{C})$: $U \in \mathcal{M}_{d'}(\mathcal{M}_d(\mathbb{C}))$. Then, the Kraus matrices L_i are (rescaled copies) of the blocks $U^{s,t} \in \mathcal{M}_d(\mathbb{C})$.

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Conclusion: almost all quantum channels are in \mathcal{C}

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From quantum channels to random density matrices

For a norm one vector $x \in \mathbb{C}^d$, define the random density matrix

$$\mathcal{U}(dd') \ni U \mapsto \rho = \Phi^{U,\beta}(|x\rangle\langle x|) = \text{Tr}_{d'}[U(|x\rangle\langle x| \otimes \beta)U^*].$$

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If $\beta = |y\rangle\langle y|$ ($y \in \mathbb{C}^{d'}$) is a rank-one projector, then

$$\rho = \text{Tr}_{d'} |U(x \otimes y)\rangle\langle U(x \otimes y)|$$

is an element of the **induced density matrices ensemble** (of parameters d, d').

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- ▶ ρ has the same distribution as $\text{Tr}_{d'} |z\rangle\langle z|$, where z is a Lebesgue-uniform vector on the unit sphere of $\mathbb{C}^d \otimes \mathbb{C}^{d'} \simeq \mathbb{C}^{dd'}$.

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- ▶ There is a connection with the **Wishart ensemble** from Random Matrix Theory: if $W = XX^*$ is a Wishart matrix of parameters d and d' , then

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- ▶ Asymptotics, in the regimes $[d \text{ fixed}, d' \rightarrow \infty]$, $[d \rightarrow \infty \text{ fixed}, d' \text{ fixed}]$ and $[d, d' \rightarrow \infty, d'/d \rightarrow c > 0]$ are well understood.

A new model of random density matrices

- ▶ induced ensemble = **one iteration** of a random channel

$$\begin{aligned}\rho &= \text{Tr}_{d'}[U(|x\rangle\langle x| \otimes |y\rangle\langle y|)U^*] \\ &= \Phi^{U,y}(|x\rangle\langle x|).\end{aligned}$$

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- ▶ We have defined an ensemble of density matrices

$$\mathcal{U}(dd') \ni U \mapsto \rho_\infty = \lim_{n \rightarrow \infty} \left(\Phi^{U,y}\right)^n(|x\rangle\langle x|).$$

The asymptotic induced ensemble

Fix $d, d' \geq 2$ and a probability vector $b \in \mathbb{C}^{d'}$. An element from the **asymptotic induced ensemble** of parameters (d, b) is the random density matrix

$$\mathcal{U}(dd') \ni U \mapsto \rho_\infty = \lim_{n \rightarrow \infty} \left(\Phi^{U, \beta} \right)^n (\rho_0),$$

where

- ▶ $\beta \in \mathcal{M}_{d'}^{1,+}(\mathbb{C})$ is a fixed density matrix with spectrum b ;
- ▶ $\rho_0 \in \mathcal{M}_d^{1,+}(\mathbb{C})$ is a fixed initial state.

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$$\mathcal{U}(dd') \ni U \mapsto \rho_\infty = \lim_{n \rightarrow \infty} \left(\Phi^{U, \beta} \right)^n (\rho_0),$$

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- ▶ The distribution of ρ_∞ is **unitarily invariant**: $\rho_\infty \stackrel{\text{law}}{=} V \rho_\infty V^*$ for all $V \in \mathcal{U}(d)$.

Numerical simulations

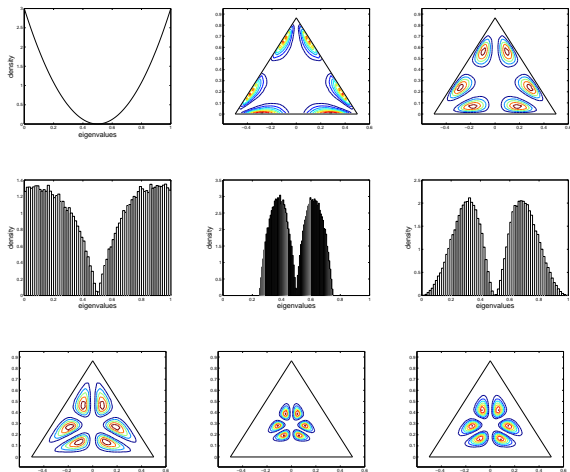


Figure: First row - induced measure $d'=2$, $d'=3$, $d'=5$; Second & third rows - asymptotic measure $b=[1, 0]$, $b=[3/4, 1/4]$, $b=[1, 0, 0, 0]$; $b=[1, 0, 0]$, $b=[3/4, 1/8, 1/8]$ and $b=[1, 0, 0, 0, 0]$.

Models of random repeated quantum interactions

Repeated quantum interactions

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We now introduce two new models:

- ▶ **random environment:** U is fixed, and the successive environment states $(\beta_n)_n$ are i.i.d. random density matrices.
- ▶ **i.i.d. unitaries:** the sequence of interaction unitaries $(U_n)_n$ is Haar-i.i.d., and no assumption is made on the (possibly random) environment states $(\beta_n)_n$.

Asymptotic results: random environment

Discrete evolution equation

$$\rho_n = \Phi^{\beta_n}(\rho_{n-1}) = \text{Tr}_{d'} [U(\rho_{n-1} \otimes \beta_n)U^*].$$

In this model, the interaction unitary U is fixed beforehand and the environment states $(\beta_n)_n$ are i.i.d. random density matrices.

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We use results by L. Bruneau, A. Joye and M. Merkli on products of random matrices, applied to the (i.i.d.) channels

$$\Phi^{\beta_n} \in \mathcal{L}(\mathcal{M}_d(\mathbb{C})).$$

Asymptotic results: random environment

Theorem (BJM)

Let $(M_n)_n$ be a sequence of i.i.d. random contractions of $\mathcal{M}_D(\mathbb{C})$ with the following properties:

1. There exists a constant vector $\psi \in \mathbb{C}^D$ such that $M\psi = \psi$ almost surely;
2. $\mathbb{P}(1 \text{ is a simple eigenvalue of } M) > 0$.

Then the (deterministic) matrix $\mathbb{E}[M]$ has eigenvalue 1 with multiplicity one and there exists a constant vector $\theta \in \mathbb{C}^D$ such that

$$\lim_{N \rightarrow \infty} \frac{1}{N} \sum_{n=1}^N M_1(\omega) M_2(\omega) \cdots M_n(\omega) = |\psi\rangle\langle\theta| = P_{1, \mathbb{E}[M]},$$

where $P_{1, \mathbb{E}[M]}$ is the rank-one spectral projector of $\mathbb{E}[M]$ corresponding to the eigenvalue 1.

Asymptotic results: random environment

Using the duality between the Schrödinger and the Heisenberg pictures of Quantum Mechanics, we obtain

Theorem

Let $(\Phi_n)_n$ be a sequence of i.i.d. random quantum channels acting on $\mathcal{M}_d(\mathbb{C})$ such that

$$\mathbb{P}(\Phi \text{ has an unique invariant state}) > 0.$$

Then $\mathbb{E}[\Phi]$ is a quantum channel with an unique invariant state $\theta \in \mathcal{M}_d^{1,+}(\mathbb{C})$ and, \mathbb{P} -almost surely,

$$\lim_{N \rightarrow \infty} \frac{1}{N} \sum_{n=1}^N [\Phi_n \circ \dots \circ \Phi_1](\rho_0) = \theta, \quad \forall \rho_0 \in \mathcal{M}_d^{1,+}(\mathbb{C}).$$

Asymptotic results: random environment

Proposition

Let $\{\beta_n\}_n$ be a sequence of i.i.d. random density matrices such that, with positive probability, the random quantum channel Φ^β has an unique invariant state. Then, almost surely, for all initial states $\rho_0 \in \mathcal{M}_d^{1,+}(\mathbb{C})$, one has

$$\lim_{N \rightarrow \infty} \frac{\rho_1 + \dots + \rho_N}{N} =$$
$$\lim_{N \rightarrow \infty} \frac{1}{N} \sum_{n=1}^N [\Phi^{\beta_n} \circ \dots \circ \Phi^{\beta_1}](\rho_0) = \theta,$$

where $\theta \in \mathcal{M}_d^{1,+}(\mathbb{C})$ is the unique invariant state of the deterministic channel $\Phi^{\mathbb{E}[\beta]}$.

In particular, if $\mathbb{E}[\beta] = I_{d'}/d'$, then θ is the “chaotic” state I_d/d .

Asymptotic results: i.i.d. unitaries

Discrete evolution equation

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In this model, the interaction unitaries U_n are Haar distributed independent random matrices.

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Lemma

Let $(V_n)_n$ be a sequence of i.i.d. Haar unitaries independent of the family $\{U_n, \beta_n\}_n$ and consider the sequence of successive states $(\rho_n)_n$ defined earlier. Then the sequences $(\rho_n)_n$ and $(V_n \rho_n V_n^)_n$ have the same distribution.*

Asymptotic results: i.i.d. unitaries

Lemma

The sequence of successive states $(\rho_n)_n$ and its i.i.d.-randomly rotated version $(V_n \rho_n V_n^)_n$ have the same distribution.*

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Proposition

Let $(\rho_n)_n$ be the successive states of a repeated quantum interaction scheme with i.i.d. random unitary interactions. Then, almost surely,

$$\lim_{n \rightarrow \infty} \frac{\rho_1 + \dots + \rho_n}{n} = \frac{I_d}{d}.$$

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- ▶ Continuous limit ?

Thank you !

<http://arxiv.org/abs/0902.2634>